

## TMA132 Fourieranalys F2/Kf2, 5 poäng

OBS! Ange namn, personnummer samt linje och inskrivningsår.

1. Ett linjärt tidsinvariant kausalt system har stegsvaret  $f(t) = e^{-t}\theta(t)$ , (dvs.  $f(t)$  är utsignalen då insignalen är  $\theta(t)$ ).

a) Beräkna utsignalen då insignalen är  $t\theta(t)$ .

b) En sinusformad insignal ger en utsignal, vars amplitud är hälften av insignalens. Beräkna vinkelfrekvensen.

2. Bestäm det polynom  $P(x)$  av högst andra graden som minimerar

$$\int_{-\infty}^{\infty} [|x| - P(x)]^2 e^{-x^2} dx.$$

3. a) Lös Laplaces differentialekvation

$$\begin{cases} u_{xx} + u_{yy} = 0, & 0 < x < 2, & y > 0, \\ u(0, y) = u_x(2, y) = 0, & \lim_{y \rightarrow \infty} u(x, y) = 0, \\ u(x, 0) = \begin{cases} 0, & 0 < x < 1, \\ 1, & 1 < x < 2. \end{cases} \end{cases}$$

b) Ge någon fysikalisk tolkning av problemet i uppgift (a).

4. Bestäm en lösning till problemet,

$$\begin{cases} u_t = u_{xx} + u_x, & -\infty < x < \infty, & t > 0, \\ u(x, 0) = x^2 e^{-x^2}, & -\infty < x < \infty. \end{cases}$$

Ledning: Fouriertransformera i  $x$ -led.

5. Låt  $f$  vara en funktion i  $L^1(\mathbb{R})$ . Definiera  $F(x) = \sum_{k=-\infty}^{\infty} f(x + 2k\pi)$ .

a) Visa att  $F$  är  $2\pi$  periodisk.

b) Härled ett samband mellan  $F$ 's Fourierkoefficienter och  $f$ 's Fouriertransform.

c) Bevisa (under lämpliga förutsättningar) *Poissons Summationsformel*:

$$\sum_{k=-\infty}^{\infty} f(2k\pi) = \frac{1}{2\pi} \sum_{n=-\infty}^{\infty} \hat{f}(n).$$

6. Bestäm temperaturen  $u = u(r, t)$  i klotet  $r = \sqrt{x^2 + y^2 + z^2} < 1$ , då

$$\begin{cases} \frac{\partial u}{\partial t} = \nabla^2 u = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}), & 0 < r < 1, & t > 0, \\ u(1, t) + u_r(1, t) = 0, & u(r, 0) = f(r), & u \text{ begränsad.} \end{cases}$$

7. Formulera och bevisa samplingsteoremet.

8.  $J_n(x)$  är Besselfunktion av ordning  $n$ . Visa genererande funktionsformel:

$$\sum_{n=-\infty}^{\infty} J_n(x) z^n = e^{\frac{x}{2}(z - \frac{1}{z})}, \quad \forall z \neq 0, \quad \forall x.$$

1. a)  $\theta(t) \xrightarrow{\mathcal{L}} \frac{1}{s+1} \xrightarrow{\text{L-transf.}} \frac{1}{s+1} = \frac{1}{s} H(s) \Rightarrow H(s) = \frac{s}{s+1}$   
 $\xrightarrow{\text{SAR}} \epsilon \theta(t) \xrightarrow{\mathcal{L}} y(t)$ . Eftersom  $\epsilon \theta(t) \xrightarrow{\mathcal{L}} \frac{1}{s^2}$  förs  
 $Y(s) = \frac{1}{s^2} H(s) = \frac{1}{s(s+1)} = \frac{1}{s} - \frac{1}{s+1}$ . Dvs

$y(t) = (1 - e^{-t}) \theta(t)$

b)  $\sin(\omega t) \xrightarrow{\mathcal{L}} A(\omega) \sin(\omega t + \varphi) = \frac{1}{2} \sin(\omega t + \varphi) \Rightarrow$   
 $|H(\omega)| = |H(i\omega)| = \left| \frac{i\omega}{i\omega+1} \right| = \frac{1}{2} \Rightarrow$

$\frac{|\omega|}{\sqrt{1+\omega^2}} = \frac{1}{2} \Rightarrow \omega^2 + 1 = 4\omega^2 \Rightarrow \omega = \frac{1}{\sqrt{3}}$

2. Använd Hermite polynomen  $H_n(x)$ . Seriv

$P(x) = a_0 H_0(x) + a_1 H_1(x) + a_2 H_2(x)$

$\int_{-\infty}^{\infty} \left[ |x| - \sum_0^2 a_n H_n(x) \right]^2 e^{-x^2} dx$  blir minimal

precis då  $a_n = c_n = \frac{1}{2^n n! \sqrt{\pi}} \int_{-\infty}^{\infty} |x| H_n(x) e^{-x^2} dx$

$H_0(x) = 1, H_1(x) = 2x, H_2(x) = 4x^2 - 2$

$c_0 = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} |x| e^{-x^2} dx = \frac{2}{\sqrt{\pi}} \int_0^{\infty} x e^{-x^2} dx = \frac{1}{\sqrt{\pi}} \left[ -e^{-x^2} \right]_0^{\infty} = \frac{1}{\sqrt{\pi}}$

$c_1 = \frac{1}{2\sqrt{\pi}} \int_{-\infty}^{\infty} |x| \cdot 2x \cdot e^{-x^2} dx = 0$  (Udda integrand)

$c_2 = \frac{1}{8\sqrt{\pi}} \int_{-\infty}^{\infty} |x| (4x^2 - 2) e^{-x^2} dx = \frac{1}{\sqrt{\pi}} \int_0^{\infty} x^3 e^{-x^2} dx - \frac{c_0}{4} =$   
 $= \frac{1}{\sqrt{\pi}} \left[ -\frac{x^2}{2} e^{-x^2} \right]_0^{\infty} + \frac{1}{\sqrt{\pi}} \int_0^{\infty} x e^{-x^2} dx = \frac{c_0}{2} - \frac{c_0}{4} = \frac{1}{4\sqrt{\pi}}$

$P(x) = \frac{1}{\sqrt{\pi}} \cdot 1 + \frac{1}{4\sqrt{\pi}} (4x^2 - 2) = \frac{1}{\sqrt{\pi}} \left( x^2 + \frac{1}{2} \right)$

3. a) (DE):  $u_{xx} + u_{yy} = 0$   $0 \leq x \leq 2, 0 \leq y < \infty$

(RV1):  $u(0, y) = 0$

(RV3):  $u(x, 0) = f(x) = \begin{cases} 0, & 0 \leq x < 1 \\ 1, & 1 \leq x < 2 \end{cases}$

(RV2):  $u'_x(2, y) = 0$

(RV4):  $\lim_{y \rightarrow \infty} u(x, y) = 0$

Delpromblem: Finn lösningar  $u(x, y) = X(x)Y(y)$  till

DE + RV1 + RV2, som inte är  $\equiv 0$ .

$X''(x)Y(y) + X(x)Y''(y) = 0; \frac{X''(x)}{X(x)} = -\frac{Y''(y)}{Y(y)} = \lambda$

Egenvärdesproblem:  $X'' = \lambda X; X(0) = X'(2) = 0$

Egentlösningar:  $X_n(x) = \sin((n + \frac{1}{2}) \frac{\pi x}{2}); n = 0, 1, 2, \dots$

Egenvärden  $\lambda_n(x) = -[(n + \frac{1}{2}) \frac{\pi}{2}]^2 = -\alpha_n^2$

$Y_n''(y) - [(n + \frac{1}{2}) \frac{\pi}{2}]^2 Y_n(y) = 0$

$Y_n(y) = A_n e^{\alpha_n y} + B_n e^{-\alpha_n y}, \alpha_n = (n + \frac{1}{2}) \frac{\pi}{2}$

$X_n(x)Y_n(y)$  löslor delpromblemet.

Superposition:

$u(x, y) = \sum_{n=0}^{\infty} [A_n e^{\alpha_n y} + B_n e^{-\alpha_n y}] \sin(\alpha_n x)$

uppfyller DE + RV1 + RV2.

RV4 uppfyllt om  $A_n = 0, \forall n$ .

$u(x, y) = \sum_{n=0}^{\infty} B_n e^{-\alpha_n y} \sin(\alpha_n x)$  löslor DE + (RV1) + (RV2) + (RV4)

(RV3):  $f(x) = \sum_{n=0}^{\infty} B_n \sin(\alpha_n x) = 0 \text{ g-serie } u(0, 2)$

$B_n = \frac{1}{H_n} \int_0^2 f(x) \sin(\alpha_n x) dx = \int_1^2 \sin(\alpha_n x) dx = [H_n = \int_0^2 \sin^2 \alpha_n x dx]$   
 $= \int_1^2 \frac{\cos((n + \frac{1}{2}) \frac{\pi x}{2})}{(n + \frac{1}{2}) \frac{\pi}{2}} dx = \frac{\cos((n + \frac{1}{2}) \frac{\pi}{2})}{(n + \frac{1}{2}) \frac{\pi}{2}}$

Svar:  $u(x, y) = \sum_{n=0}^{\infty} \frac{\cos \alpha_n}{\alpha_n} e^{-\alpha_n y} \sin(\alpha_n x), \alpha_n = (n + \frac{1}{2}) \frac{\pi}{2}$

b) Värmeledning i området  $0 < x < 2, 0 < y < \infty, -\infty < z < \infty$ .

Begränsningsytan  $x=0$  hålls vid temp  $u=0$ , ytan  $x=2$  är isolerad, ytan  $y=0$  vid temp.  $f(x)$ . Stationär tillstånd. Inga inre värme källor.

$$4. \begin{cases} u_t = u_{xx} + u_x, & -\infty < x < \infty, t > 0 \\ u(x, 0) = x^2 e^{-x^2} = f(x). \end{cases}$$

Fouriertransformera i  $x$ -led.  $\mathcal{F}_x [u(x, t)] = \hat{u}(\xi, t)$ .

$$\frac{\partial \hat{u}}{\partial t} = (i\xi)^2 \hat{u} + (i\xi) \hat{u} = (-\xi^2 + i\xi) \hat{u}$$

$$\hat{u}(\xi, t) = C(\xi) e^{(-\xi^2 + i\xi)t}$$

$$\hat{u}(\xi, 0) = C(\xi) = \hat{f}(\xi)$$

$$e^{-x^2} \mathcal{F} \Rightarrow \frac{1}{\sqrt{\pi}} e^{-\xi^2/4} \Rightarrow x e^{-x^2} \mathcal{F} \Rightarrow \frac{d}{d\xi} \left( \frac{1}{\sqrt{\pi}} e^{-\xi^2/4} \right) = -\frac{i\sqrt{\pi}}{2} \xi e^{-\xi^2/4}$$

$$f(x) = x^2 e^{-x^2} \mathcal{F} \Rightarrow \frac{d}{d\xi} \left( -\frac{i\sqrt{\pi}}{2} \xi e^{-\xi^2/4} \right) = \frac{\sqrt{\pi}}{2} \left( 1 - \frac{\xi^2}{2} \right) e^{-\xi^2/4} = \hat{f}(\xi)$$

$$\hat{u}(\xi, t) = \frac{\sqrt{\pi}}{2} \left( 1 - \frac{\xi^2}{2} \right) e^{-\xi^2/4} e^{-\xi^2 t} e^{i\xi t} =$$

$$= \frac{\sqrt{\pi}}{2} \left( 1 - \frac{\xi^2}{2} \right) e^{-(t + 1/4)\xi^2} e^{i\xi t}$$

Lat  $A = t + 1/4$ ;  $\frac{1}{\sqrt{4\pi A}} e^{-x^2/4A} \mathcal{F} e^{-A\xi^2}$

$$-i \frac{d}{dx} \left( \frac{1}{\sqrt{4\pi A}} e^{-x^2/4A} \right) \mathcal{F} \xi e^{-A\xi^2}$$

$$\xi e^{-A\xi^2} \subset \frac{1}{\sqrt{4\pi A}} \frac{ix}{2A} e^{-x^2/4A}$$

$$\xi^2 e^{-A\xi^2} \subset -i \frac{d}{dx} \left( \frac{1}{\sqrt{4\pi A}} \frac{ix}{2A} e^{-x^2/4A} \right) = \frac{1}{2A\sqrt{4\pi A}} \left( 1 - \frac{x^2}{2A} \right) e^{-x^2/4A}$$

$$\frac{\sqrt{\pi}}{2} \left( 1 - \frac{\xi^2}{2} \right) e^{-A\xi^2} \subset \frac{\sqrt{\pi}}{2} \frac{1}{\sqrt{4\pi A}} \left[ 1 - \frac{1}{2} \cdot \frac{1}{2A} \left( 1 - \frac{x^2}{2A} \right) \right] e^{-x^2/4A}$$

$$= \frac{1}{2} \cdot \frac{1}{\sqrt{4t+1}} \left[ 1 - \frac{1}{4t+1} \left( 1 - \frac{x^2}{4t+1} \right) \right] e^{-x^2/4t+1}$$

$$= \frac{2t(4t+1) + (x+t)^2}{(4t+1)^{5/2}} e^{-x^2/4t+1}$$

$$u(x, t) = \frac{2t(4t+1) + (x+t)^2}{(4t+1)^{5/2}} e^{-\frac{(x+t)^2}{4t+1}}$$

5. a)  $F(x) = \sum_{-\infty}^{\infty} f(x+2k\pi)$

$F(x+2\pi) = \sum_{-\infty}^{\infty} f(x+2(k+1)\pi) = \sum_{-\infty}^{\infty} f(x+2k\pi) = F(x)$

$\therefore F$   $2\pi$ -periodisk.

b)  $C_n(F) = \frac{1}{2\pi} \int_{-\pi}^{\pi} F(x) e^{-inx} dx = \frac{1}{2\pi} \sum_{-\infty}^{\infty} \int_{-\pi}^{\pi} f(x+2k\pi) e^{-inx} dx$   
 $= \sum_{-\infty}^{\infty} \frac{1}{2\pi} \int_{(2k-1)\pi}^{(2k+1)\pi} f(x) e^{-inx} dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-inx} dx = \hat{f}(n)$

c)  $\sum_{-\infty}^{\infty} f(2k\pi) = F(0) = \sum_{-\infty}^{\infty} C_n(F) e^{in \cdot 0} = \frac{1}{2\pi} \sum_{-\infty}^{\infty} \hat{f}(n)$

6.  $\begin{cases} \frac{u}{rt} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}); & 0 < r < 1; t > 0, \\ u(r,t) \text{ begränsad} \\ u(r,t) + u_r(r,t) = 0 \\ u(r,0) = f(r) \end{cases}$

Vi kan skriva  $\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \frac{\partial u}{\partial r}) = \frac{1}{r} \frac{\partial^2}{\partial r^2} (ru)$

För  $v = ru$  får då ekvationerna:

$\begin{cases} \frac{v}{\partial t} = \frac{\partial^2 v}{\partial r^2} \\ v(0,t) = v_r(1,t) = 0, \quad v(r,0) = r f(r) \end{cases}$

$v(r,t) = R(r)T(t)$  i de homogena ekvationer ger  $\frac{T'}{T} = \frac{R''}{R} = -\lambda$

$R'' + \lambda R = 0, R(0) = R'(1) = 0 \Rightarrow R = R_n(r) = \sin((n+\frac{1}{2})\pi r)$

$T' = -\lambda_n T \Rightarrow T = T_n(t) = C_n e^{-\lambda_n t}$  Ansatz

$v(r,t) = \sum_0^{\infty} C_n e^{-\lambda_n t} \sin((n+\frac{1}{2})\pi r) \Rightarrow v(r,0) = \sum_0^{\infty} C_n \sin((n+\frac{1}{2})\pi r) = r f(r)$

$\{\sin((n+\frac{1}{2})\pi r)\}_0^1$  orthonormal bas  $\Rightarrow C_n = \frac{2 \int_0^1 r f(r) \sin((n+\frac{1}{2})\pi r) dr}{\int_0^1 \sin^2((n+\frac{1}{2})\pi r) dr}$

Med denna  $C_n$  blir lösningen  $u(r,t) = \frac{1}{r} \sum_{n=0}^{\infty} C_n e^{-(n+\frac{1}{2})^2 \pi^2 t} \sin((n+\frac{1}{2})\pi r)$   
 Ans.