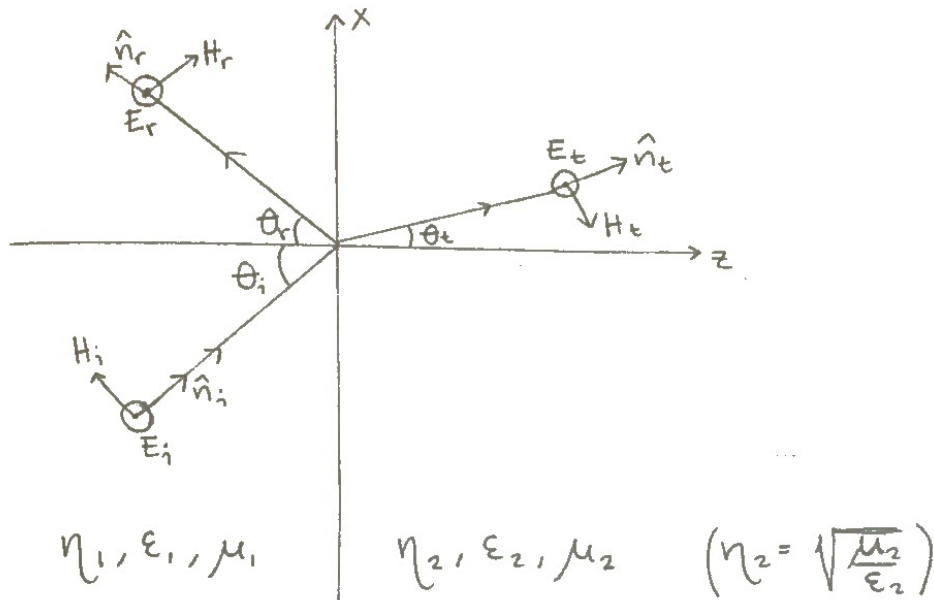


Storgruppsövning 10/12-13

Oblique incidence at a plane dielectric boundary.

Perpendicular polarization

E-field is perpendicular to the plane of incidence.

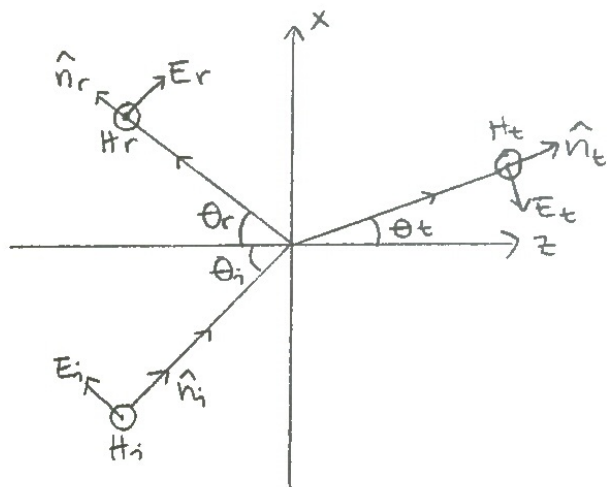


$$\begin{cases} \Gamma_{\perp} = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 \cos \theta_i - \eta_1 \cos \theta_t}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t} \\ \mathcal{T}_{\perp} = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2 \cos \theta_i}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t} \end{cases}$$

$$1 + \Gamma_{\perp} = \mathcal{T}_{\perp}$$

Parallel polarization

E-field is lying in the plane of incidence.



$$\begin{cases} \Gamma_{\parallel} = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i} \\ \mathcal{T}_{\parallel} = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2 \cos \theta_i}{\eta_2 \cos \theta_i + \eta_1 \cos \theta_t} \end{cases}$$

$$1 + \Gamma_{\parallel} = \mathcal{T}_{\parallel}$$

No reflection $\iff \theta_i = \text{Brewster angle} = \theta_B$

$$\Gamma_{\perp} = 0 \implies \eta_2 \cos \theta_i = \eta_2 \cos \theta_t$$

$$\cos \theta_t = \sqrt{1 - \sin^2 \theta_t} \xrightarrow{\text{Snell's law}} \cos \theta_t = \sqrt{1 - \left(\frac{\eta_1}{\eta_2}\right)^2 \sin^2 \theta_t}$$

$$\implies \sin \theta_i = \sin \theta_{B\perp} = \frac{1 - \mu_1 \epsilon_2 / \mu_2 \epsilon_1}{1 - \left(\frac{\mu_1}{\mu_2}\right)^2}$$

if $\mu_1 = \mu_2 \implies \theta_B$ does not exist.

$$\Gamma_{\parallel} = 0 \implies \begin{cases} \eta_2 \cos \theta_t = \eta_1 \cos \theta_i \\ \cos \theta_t = \sqrt{1 - \sin^2 \theta_t} \end{cases} \implies \sin \theta_i = \sin \theta_{B\parallel} = \frac{1 - \mu_2 \epsilon_1 / \mu_1 \epsilon_2}{1 - \left(\frac{\epsilon_1}{\epsilon_2}\right)^2}$$

$$\mu_2 = \mu_1 \implies \sin \theta_{B\parallel} = \frac{1}{\sqrt{1 + \frac{\epsilon_1}{\epsilon_2}}} \implies \tan \theta_{B\parallel} = \frac{\eta_2}{\eta_1}$$

12.12

Circular cross section wire.

Radius $a = 0,1 \text{ mm}$

$\sigma = 5 \cdot 10^6 \text{ S/m}$

$\mu_r = 100$.

Calculate the ratio of the resistance at $f_1 = 50 \text{ Hz}$ & $f_2 = 10 \text{ MHz}$

Skin depth for good conductor, $\delta = \sqrt{\frac{1}{\pi f \mu \sigma}}$ ($\sigma \omega \epsilon \gg 1$)

$$\begin{cases} f_1 = 50 \text{ Hz} \implies \delta_1 = 3,2 \cdot 10^{-3} \\ f_2 = 10 \cdot 10^6 \implies \delta_2 = 7 \cdot 11 \cdot 10^{-6} \end{cases} \quad (a = 0,1 \text{ mm})$$

$\delta_1 \gg a \implies$ current dist. on the whole cross-section.

$\delta_2 \ll a \implies$ —||— on a thin layer

$$R_1 = \frac{l}{\sigma S_1} = \frac{l}{\sigma \pi a^2}$$

$$R_2 = \frac{l}{\sigma S_2} = \frac{l}{\sigma 2\pi a \delta_2}$$

$$\frac{R_1}{R_2} = \frac{S_2}{S_1} = \frac{2\delta_2}{a} \approx 0,142 \dots$$

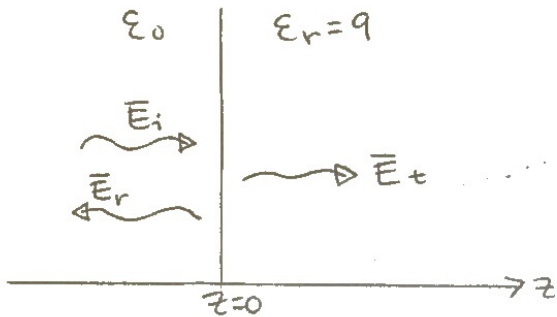
The resistance at 50 Hz is approx 14% of that for 10 MHz.

13.7

A plane wave in vacuum incident normally to a flat surface at $z=0$, of a lossless dielectric, $\epsilon_r=9$.

$$\vec{E} = \hat{x} 10 \cos(\omega t - \beta z), \quad f = 300 \text{ MHz}$$

Find the location of max for E-field in vacuum,
 z max(E_{tot}),



$$\vec{E}_i = \hat{x} 10 e^{-i\beta z}, \quad \vec{E}_r = \nabla \vec{E}_i = \nabla \hat{x} 10 e^{+i\beta z}$$

$$\vec{E}_{tot} = \vec{E}_i + \vec{E}_r$$

$$\Gamma = \text{refl. coef.} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = \frac{\sqrt{\frac{\mu_0}{\epsilon_0 \epsilon_r}} - \sqrt{\frac{\mu_0}{\epsilon_0}}}{\sqrt{\frac{\mu_0}{\epsilon_0 \epsilon_r}} + \sqrt{\frac{\mu_0}{\epsilon_0}}} = \frac{1/\sqrt{\epsilon_r} - 1}{1/\sqrt{\epsilon_r} + 1} = \frac{1/3 - 1}{1/3 + 1} = -\frac{1}{2}$$

$$\Rightarrow \vec{E}_r = -5 \hat{x} e^{+i\beta z}$$

$$\vec{E}_{tot,1} = \hat{x} (10 e^{-i\beta z} - 5 e^{i\beta z}) = \hat{x} (10 - 5 e^{2i\beta z}) e^{-i\beta z}$$

$$\Rightarrow |\vec{E}_{tot,1}| = |10 - 5 e^{2i\beta z}| = \left| \underbrace{10 - 5 \cos(2\beta z)}_{\text{Re}} - \underbrace{5 i \sin(2\beta z)}_{\text{Im}} \right|$$

$$|\vec{E}_{tot,1}| = \sqrt{(10 - 5 \cos(2\beta z))^2 + 25 \sin^2(2\beta z)} = \sqrt{25 - 100 \cos(2\beta z)}$$

$-1 \Rightarrow \text{max}$

$$\cos(2\beta z) = -1 \Rightarrow |\vec{E}_{tot,1}|_{\text{max}} = \sqrt{225} = 15$$

$$2\beta z_{\text{max}} = -(2n+1)\pi \Rightarrow z_{\text{max}} = \frac{-(2n+1)\pi}{2\beta}$$

$$\beta = \frac{\omega}{c} = \frac{2\pi \cdot 3 \cdot 10^8}{3 \cdot 10^8} = 2\pi \Rightarrow z_{\text{max}} = \frac{-(2n+1)\pi}{2 \cdot 2\pi} = \frac{-(2n+1)}{4}$$

$(n=0, 1, 2, \dots)$

13.14

A light beam is broken and totally reflected in a lossless prism. Refraction is at Brewster angle. Return wave is parallel to incident wave. Find a suitable range for (n) .

① Brewster angle at 1st interface:

$$\tan \theta_B = \frac{n}{n_0} = n$$

$$\tan(\pi/2 - \alpha) = n$$

② Snell's law of refr., 1st interf.:

$$\sin(\pi/2 - \alpha) = n \sin \beta$$

③ Total reflection, 2nd interf.:

$$\gamma > \theta_c$$

④ Snell's law, 2nd interf.:

$$\sin \theta_c \cdot n = \sin \pi/2 \cdot 1$$

$$\sin \theta_c = 1/n$$

⑤ In triangle:

$$\alpha + (\pi/2 + \beta) + \gamma = \pi$$

Suppose $\gamma = \theta_c$ ($\tan \alpha = 1/n$)

① $\sin(\pi/2 - \alpha) = n \cos(\pi/2 - \alpha) \Rightarrow \cos \alpha = n \sin \alpha$

② $\cos \alpha = n \sin \beta$
 $\Rightarrow \alpha = \beta$

④, ⑤: $\sin \gamma = \sin \theta_c = \sin(\pi/2 - (\alpha + \beta)) = 1/n$

$$\Rightarrow \cos(\alpha + \beta) = 1/n \Rightarrow \cos 2\alpha = 1/n$$

$$\cos(2\alpha) = \frac{1 - \tan^2 \alpha}{1 + \tan^2 \alpha} = 1/n \Rightarrow \{\tan(\alpha) = 1/n\} \Rightarrow \frac{1 - (1/n)^2}{1 + (1/n)^2} = 1/n$$

$$\Rightarrow n^3 - n^2 - n - 1 = 0$$

$$\Rightarrow n_c \approx 1,839$$

If n increase \Rightarrow from ① $\alpha = \beta$ will decrease $\Rightarrow \gamma$ will decrease
 from ④ θ_c will decrease

So if $n \geq n_c \Rightarrow \gamma \geq \theta_c \Rightarrow$ total refraction at 2nd interface

